

# Energy Fluxes in RMHD Turbulence

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Plasma is everywhere!

- Interstellar medium, Stars, Lab plasmas.

Important applications:

- Solar/Stellar activity, Space weather, Fusion energy.
- We want to model it as accurately as possible.
- Ideally use kinetic theory. . . but too many particles to follow. . .
- $\Rightarrow$  Fluid approximation: Compressible MHD model
- Widely applicable, but. . . very demanding. . .
- need more assumptions. . .
- but more assumptions means less applicable. . .
- **Research Question:** How can we identify which dynamics are captured and which are not?

# More assumptions...

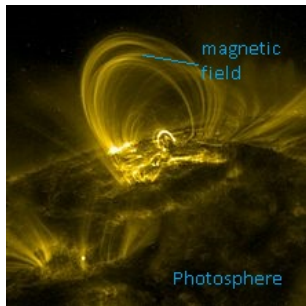


Image  
of coronal loops taken from  
<https://scied.ucar.edu/sun-coronal-loops> from  
TRACE/NASA with added  
text.

In many applications there is/are:

- Dominant magnetic field direction,
- Small  $\parallel$  components,
- Small  $\parallel$  gradients.

$\Rightarrow$  reduced MHD (RMHD) model...

- Origins in fusion
- What processes are captured?
- Can it be used in other fields?
- Coronal loops?
- Turbulence?

# What is RMHD? (Kadomtsev74, Strauss76, Zank93)[3, 4, 5]

A non-linear, low frequency approximation of 3D MHD.

$$\underline{B} = \frac{B_0}{\varepsilon_{KP}} \hat{e}_z + \underline{b}^{(1)} + \varepsilon_{KP} \underline{b}^{(2)} \dots, \quad \underline{v} = \underline{v}^{(1)} + \varepsilon_{KP} \underline{v}^{(2)} \dots,$$

Assumptions:

①  $B_0$  'large',  
 $\Rightarrow \varepsilon_{KP} = b^2/B_0^2 \ll 1$ .

② spectral anisotropy

$$\Rightarrow \varepsilon_S = k_{\parallel} \ll k_{\perp},$$

③ variance anisotropy

$$\beta \lesssim 1:$$

$$\Rightarrow \underline{B}_0 \cdot \underline{v} = \underline{B}_0 \cdot \underline{b} = 0,$$

- $\nabla_{\perp} \cdot \underline{b} = 0, \quad \nabla_{\perp} \cdot \underline{v} = 0$ .
- **incompressible**  $\perp$  only ( $\partial_z$  small).
- similar to (but not quite!) 2D MHD.

$$\text{If } \beta \lesssim 1: b_z^{(1)} = v_z^{(1)} = 0,$$

If  $\beta \gg 1: b_z^{(1)}, v_z^{(1)} \neq 0$ , passive scalars

The RMHD equations (leading order):

$$\partial_t \underline{v} + \underline{v} \cdot \nabla_{\perp} \underline{v} = \underline{b} \cdot \nabla_{\perp} \underline{b} + B_0 \partial_z \underline{b} - \nabla_{\perp} P + \nu \nabla_{\perp}^2 \underline{v},$$

$$\partial_t \underline{b} + \underline{v} \cdot \nabla_{\perp} \underline{b} = \underline{b} \cdot \nabla_{\perp} \underline{v} + B_0 \partial_z \underline{v} + \eta \nabla_{\perp}^2 \underline{b}.$$

# Dynamics captured in RMHD?

- **Problem:** When is RMHD applicable?
- **Aim:** Identify processes included/neglected in RMHD.
- **Method:** Apply recent framework to RMHD and compare.
- HD framework: Johnson21 [2], MHD: Capocci25 [1].
- Splits energy flux into physically meaningful measurable terms:
  - Vortex stretching.
  - Strain rate self-amplification.
  - Current filament stretching.
  - Current sheet thinning.
- Process:
  - Gaussian filter: Remove small scales below  $\ell \Rightarrow$  Focus on inertial scales.
  - Introduce subgrid-scale (SGS) stress tensor:
$$\tau^\ell(f_i, g_j) = \overline{f_i g_j}^\ell - \overline{f_i}^\ell \overline{g_j}^\ell.$$
  - $\tau \Rightarrow$  function of field gradients.
  - Field gradients  $\Rightarrow S \Omega J \Sigma$ .

# Incompressible MHD

$$\partial_t u_i + \partial_j(u_i u_j) = -\partial_i P + \partial_j(b_i b_j),$$

$$\partial_t b_i + \partial_j(b_i u_j) = \partial_j(u_i b_j),$$

$$\partial_i u_i = 0, \quad \partial_i b_i = 0.$$

$$\partial_t E_v + \nabla \cdot \mathcal{J}_v = -\Pi^{I,\ell} - \Pi^{M,\ell} - W^\ell,$$

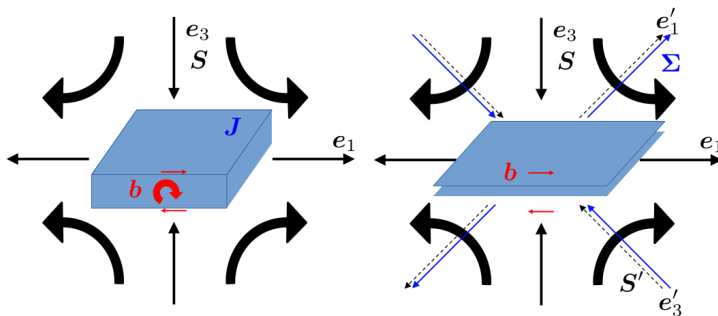
$$\partial_t E_b + \nabla \cdot \mathcal{J}_b = -\Pi^{A,\ell} - \Pi^{D,\ell} + W^\ell$$

- Inertial:  $\Pi^{I,\ell} = -\partial_j \bar{v}_i^\ell \tau^\ell(v_i, v_j)$ .
- Maxwell:  $\Pi^{M,\ell} = \partial_j \bar{v}_i^\ell \tau^\ell(b_i, b_j)$ .
- Advection:  $\Pi^{A,\ell} = -\partial_j \bar{b}_i^\ell \tau^\ell(b_i, v_j)$ .
- Dynamo:  $\Pi^{D,\ell} = \partial_j \bar{b}_i^\ell \tau^\ell(v_i, b_j)$ .
- Energy Conversion:  $W^\ell = \bar{b}_i^\ell \bar{b}_j^\ell \partial_j \bar{v}_i^\ell$ .
- $> 0 \Rightarrow$  forward cascade.
- $\Pi$  are further split into subfluxes

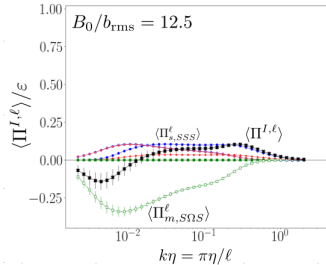
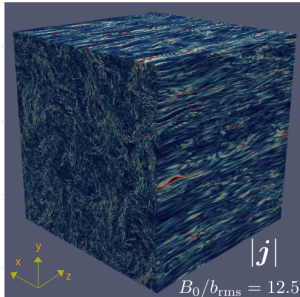
# MHD Dynamics (Capocci2025)

- Subfluxes are triple product of combinations of  $S$ ,  $\Omega$ ,  $J$ ,  $\Sigma$ .
- Example: Current sheet thinning -  

$$\Pi_{s, SJ\Sigma}^{M, \ell} \propto \text{Tr} \left( \bar{S}^{\ell} \left( \bar{J}^{\ell} \bar{\Sigma}^{\ell} - \bar{\Sigma}^{\ell} \bar{J}^{\ell} \right) \right).$$



# MHD Results (Capocci2025): Inertial Flux $\Pi'$



vortex thinning

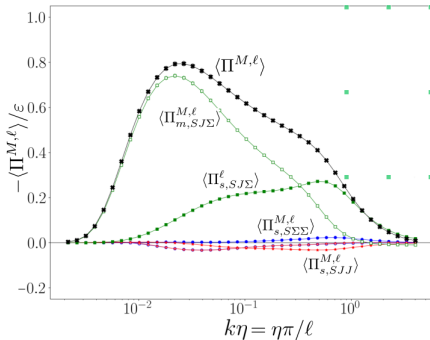
$$\Pi_{m,\Omega S}^\ell = - \int_0^{\ell^2} d\theta \operatorname{tr} \left( \overline{S}^\ell \left( \overline{S^{\sqrt{\theta}} \Omega^{\sqrt{\theta}} \phi} - \overline{\Omega^{\sqrt{\theta}} S^{\sqrt{\theta}} \phi} \right) \right)$$

⇒ Main contribution: Vortex thinning.



# MHD Results (Capocci2025): $\Pi^M$

$$\Pi^{M,\ell} = \underbrace{\Pi_{s,S\Sigma\Sigma}^{M,\ell} + \Pi_{m,S\Sigma\Sigma}^{M,\ell}}_{\text{extensional/restoring effect of magnetic field}} + \underbrace{\Pi_{s,SJJ}^{M,\ell} + \Pi_{m,SJJ}^{M,\ell}}_{\text{strain amplification by current-filament stretching}} + \underbrace{\Pi_{s,SJ\Sigma}^{M,\ell} + \Pi_{m,SJ\Sigma}^{M,\ell}}_{\text{small-scale strain amplification by current-sheet thinning}}$$



$$\bar{\Sigma}^\ell = (\nabla \bar{\mathbf{b}}^\ell + (\nabla \bar{\mathbf{b}}^\ell)^\ell) / 2$$

$$\bar{\mathcal{J}}^\ell = (\nabla \bar{\mathbf{b}}^\ell - (\nabla \bar{\mathbf{b}}^\ell)^\ell) / 2$$

Energy transfer almost exclusively

$$\Pi_{m,SJ\Sigma}^{M,\ell} =$$

$$-\int_0^{\ell^2} d\theta \operatorname{tr} \left( \bar{\mathcal{S}}^\ell \left( \overline{\mathcal{J}^{\sqrt{\theta}} \Sigma^{\sqrt{\theta}}}^\phi - \overline{\Sigma^{\sqrt{\theta}} \mathcal{J}^{\sqrt{\theta}}}^\phi \right) \right)$$

$$\Pi_{s,SJ\Sigma}^{M,\ell} \approx -\ell^2 \operatorname{tr} \left( \bar{\mathcal{S}}^\ell (\bar{\mathcal{J}}^\ell \bar{\Sigma}^\ell - \bar{\Sigma}^\ell \bar{\mathcal{J}}^\ell) \right)$$

⇒ Main contribution: Current sheet thinning and flow response.

## Current work: Apply this to RMHD

The leading order RMHD equations:

$$\partial_t v_i + \partial_j(v_i v_j) = -\partial_{i, \neq z}(P) + \partial_j(b_i b_j) + B_0 \partial_z b_i,$$

$$\partial_t b_i + \partial_j(b_i v_j) = \partial_j(v_i b_j) + B_0 \partial_z v_i.$$

$$\partial_j \cdot v_i = \partial_j b_i = 0, \quad j = x, y.$$

$$\Rightarrow \partial_t E_v + \nabla \cdot \mathcal{J}_v = -\Pi^{I,\ell} - \Pi^{M,\ell} - W_{\text{RMHD}}^\ell - B_0 \partial_z \bar{v}_i^\ell \bar{b}_i^\ell,$$

$$\partial E_b + \nabla \cdot \mathcal{J}_b = -\Pi^{A,\ell} - \Pi^{D,\ell} + W_{\text{RMHD}}^\ell + B_0 \partial_z \bar{v}_i^\ell \bar{b}_i^\ell$$

- Same definitions, except  $j \neq z$ .
- $W_{\text{RMHD}}^\ell$  only contains  $\perp$  gradients.
- $B_0 \partial_z$  is the only  $z$  derivative:  $dz$  small,  $B_0$  large.
- In theory expect same dominant processes in RMHD.
- Vortex thinning more important as  $B_0$  increases.

- Compare strain rate tensor for MHD and RMHD:

$$S_{\text{MHD}} = \begin{bmatrix} \partial_x v_x & \frac{1}{2}(\partial_y v_x + \partial_x v_y) & \frac{1}{2}(\partial_z v_x + \partial_x v_z) \\ \frac{1}{2}(\partial_x v_y + \partial_y v_x) & \partial_y v_y & \frac{1}{2}(\partial_z v_y + \partial_y v_z) \\ \frac{1}{2}(\partial_x v_z + \partial_z v_x) & \frac{1}{2}(\partial_y v_z + \partial_z v_y) & \partial_z v_z \end{bmatrix},$$
$$S_{\text{RMHD}} = \begin{bmatrix} \partial_x v_x & \frac{1}{2}(\partial_y v_x + \partial_x v_y) & \frac{1}{2}\partial_x v_z \\ \frac{1}{2}(\partial_x v_y + \partial_y v_x) & \partial_y v_y & \frac{1}{2}\partial_y v_z \\ \frac{1}{2}\partial_x v_z & \frac{1}{2}\partial_y v_z & 0 \end{bmatrix}.$$

- Low- $\beta \sim 2\text{D}$ 
  - Dominated by vortex and current sheet thinning
- Terms involving  $\parallel$  components  $\Rightarrow$  high- $\beta$ .

# Summary

Identify dynamics that are captured in RMHD.

- Low  $\beta$ :
  - $\sim$  2D as in Capocci2025.
  - Nonzero subfluxes:
    - Vortex thinning:  $\Pi_{s, s\Omega s}^{I, \ell}$
    - Current sheet thinning and flow response:  $\Pi_{sJ\Sigma}^{M, \ell}$
  - Also dominant in full 3D MHD using DNS.
    - Vortex thinning becomes stronger as  $B_0$  increases.
- Current work  $\Rightarrow$  High  $\beta$ :
  - Expect current filament stretching to become more important.
  - Understand how the  $\parallel$  components affect the dynamics.

Thank you!

Questions? Comments? Feedback?

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